

Outline

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TI DOE

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Difference Imaging

Bonus! Scattered/Extraneous Light

What's different about these chips?

The main effects that we're now facing are:

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- Thick devices mean that the PSF is a function of the conversion depth This leads to interesting cosmic rays and chromatic terms
- There is a lateral redistribution of charge This leads to noise correlations and non-Poisson statistics
- Non-uniformities in the Si affect the pixel size This leads to astrometric, photometric, and shape anomalies

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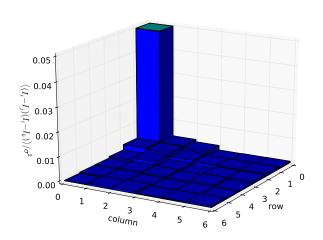
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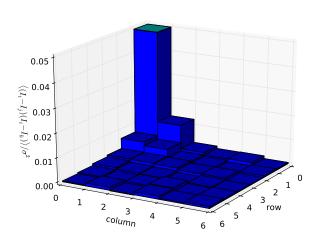
Correlations

Visits 124210, 124212, CCDs 0-9 $\langle I \rangle$ = 2392



Correlations

Visits 127566, 127567, CCDs 0-9 $\langle I \rangle = 28241$



If each photon is detected independently, we have

$$Var(I) = I/g$$

where g is the gain (i.e. the number of ADU per electron).

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where g is the gain (i.e. the number of ADU per electron). If up-fluctuations decrease the probability of collecting another charge, the variance is decreased, and the apparent gain increases with the charge level (and thus variance)

Level		1850	23500	32500	65000
CCD	Amp	gain	gain	gain	gain
I_54	I	2.87	3.19	3.21	3.35
I_54	2	3.02	3.20	3.22	3.34
I_54	3	2.93	3.14	3.15	3.26
I_54	4	2.99	3.21	3.24	3.35
I_55	1	3.68	3.92	3.94	4.08
I_55	2	3.52	3.74	3.78	3.91
I_55	3	3.44	3.70	3.72	3.86
I_55	4	3.46	3.69	3.73	3.87
I_56	1	3.15	3.35	3.37	3.46
I_56	2	3.16	3.34	3.38	3.48
I_56	3	3.07	3.26	3.29	3.40
I_56	4	3.14	3.31	3.35	3.45

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The measured non-linearity is about 1%; we are seeing the effect of charge migration

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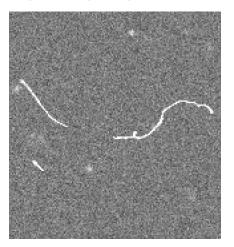
Astrometry

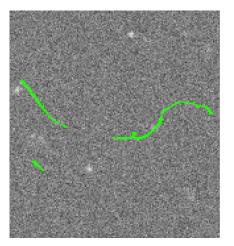
The PS

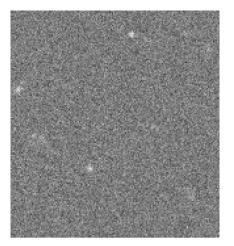
Shapes

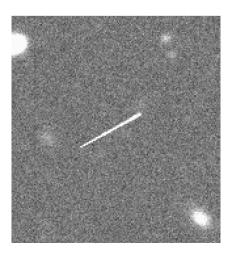
Difference Imaging

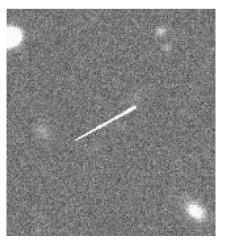
Bonus! Scattered/Extraneous Light



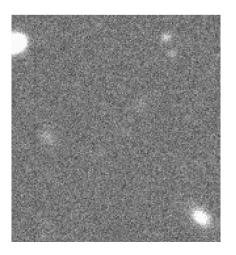








Not only are the tracks long, but they're fatter at one end than the other



Cosmic Ray Removal

Perhaps surprisingly, the algorithm we used in SDSS works here too.

Cosmic Ray Removal

Cosmic ray contaminated pixels satisfy a series of conditions:

- 1. That the candidate bad pixel p not be adjacent to a saturated pixel.
- 2. That p's intensity I exceed the local background by $n\sigma(B)$ ($\sigma(I)$ is the standard deviation of I; $n \approx 6$).
- 3. That no pixel be part of a peak which is sharper than the centre of a star centred in a pixel; i.e.

$$I - c\sigma(I) > c_2\phi(d)\left(\overline{I} + c\sigma(\overline{I})\right)$$

where c and c_2 are constants (\approx 3.0 and 0.6), $\phi(d)$ is the value of the PSF at a distance d from the centre of a star, and \bar{l} is the average of two pixels a distance d away from p.

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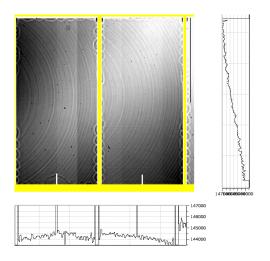
The PS

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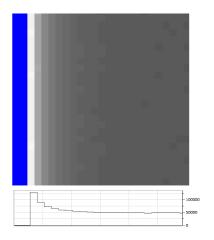
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Tree Rings

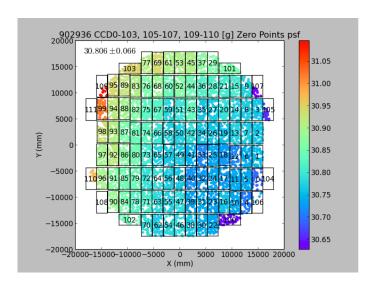


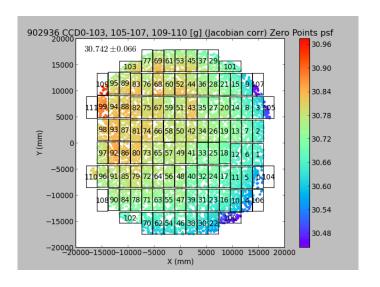
HSC ``Tree rings" at the 0.3% level

Pixel Size Variation at the Device Edges

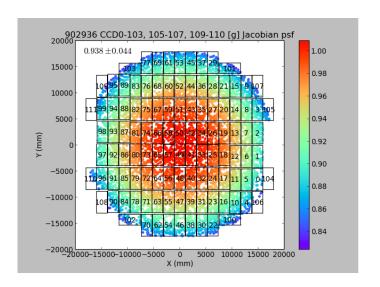


The electric field diverges near the edge of the CCD, and this leads to larger pixels (by a factor of c. 200% at the very edge of an HSC device).





This isn't a CCD effect; it's due to the distortion in the optics.



This is a variation of the size of the pixels, and can be handled in the same way

This is a variation of the size of the pixels, and can be handled in the same way ... providing the tree rings have a scale larger compared to sources.

Variable Pixel Sizes

Variable pixel sizes have a different effects on different measures of flux (e.g. aperture and psf), but it appears likely that this can be corrected at the catalogue level.

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Variable pixel sizes have a different effects on different measures of flux (e.g. aperture and psf), but it appears likely that this can be corrected at the catalogue level.

A nasty (theoretical?) problem is that if the pixels are individuals then we become sensitive to the pixel substructure of the pixels even if we are well sampled.

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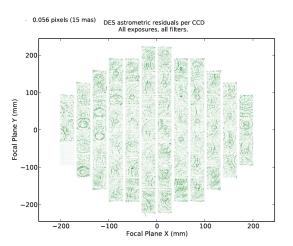
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Tree Rings in DES



DES sees 2-3% treerings, and they affect the astrometry (at the few mas level?). We don't expect significant problems in HSC.

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wcs = data.getWcs()
print wcs.pixelToSky(afwGeom.PointD(10, 100))
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In other words, that astrometry is a mapping from the pixel coordinates to the sky --- note that I didn't write ``pixel grid". There's another mapping that we need to carry in the software.

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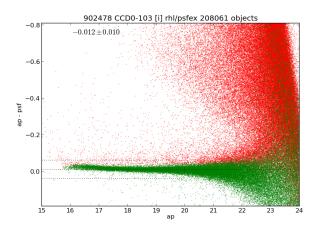
Intensity-Dependent PSFs

We may compare the aperture and PSF-model measurements of stellar fluxes. Redistributing charge within the aperture has no effect on the total flux, whereas PSF measures are sensitive to the star's profile.

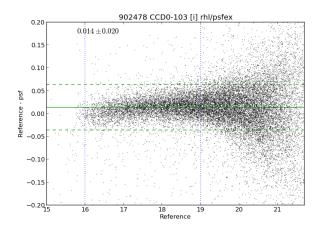
Intensity-Dependent PSFs

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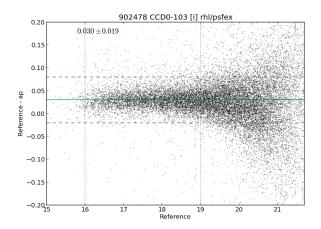
PSF - Aperture Magnitudes



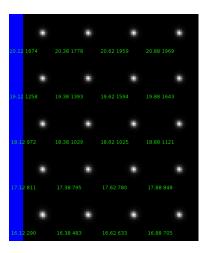
PSF v. SDSS



PSF v. SDSS

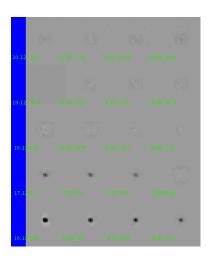


The PSF



Average star images in 0.25 mag wide bins (16--16.25 mag in bottom left)

The PSF



The residuals resulting from subtracting the average star in the 19--19.25 mag bin from all the average stars in the previous mosaic. The core of the bottom left star corresponds to $\sigma \sim 0.75$ pixels.

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$$\delta^{X}_{ij} = a^{X}_{ij} Q_{ij}.$$

The charge density at this boundary is $({\bf Q}_{00}+{\bf Q}_{\rm X})/2$, so the charge in the charge in pixel (0,0) is

$$\delta Q_{00} \propto \sum_{ij} \sum_{X} \delta_{ij}^{X} (Q_{00} + Q_{X})/2 \sim \sum_{ij} \sum_{X} a_{ij}^{X} Q_{ij} (Q_{00} + Q_{X})$$

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- the correlation function of dome flats
- the as-observed image from any intensity distribution

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The Pierres tells me that, with some assumptions, the correlation function may be used to estimate the a coefficients.

Modelling the PSF

If we have such a model, we can parameterize the zero-intensity PSF and create the as-observed image of any star in the field; this permits us to estimate the zero-intensity PSF.

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This makes the mathematics a bit nastier, but we need to do the work anyway to handle:

- undersampling
- distortions such as tree-rings
- optical aberrations (e.g. coma)

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PSF effects on Shape Measurement

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We know how to allow for the intensity-dependence of the PSF; if we're dealing with faint sources we are done --- but theoretically we can apply the same model to the data. This can also be used for coadd-based approaches using a per-object coadd PSF.

Variable Pixel Size

Do something! If we can measure this we can correct for it...

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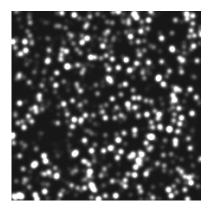
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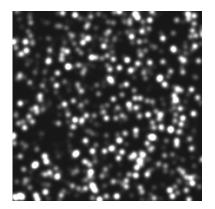
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Template

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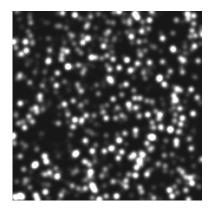
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Data

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Matched Template

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Difference

This intensity-dependence of the PSF (and the shapes of galaxies) has deleterious effects upon image subtraction.

For example, if the data exposure has half the count level of the template and 10% worse seeing:

We have to correct both the template and data for this effect before building the difference kernel

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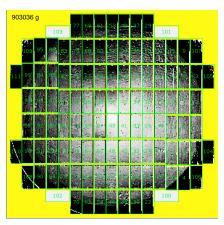
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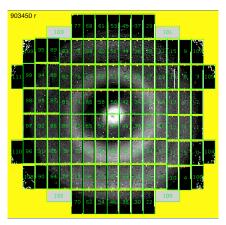
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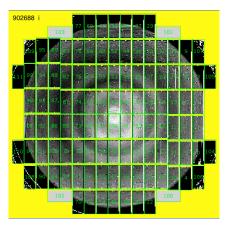
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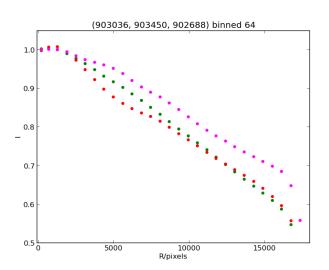


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The End

Perhaps surprisingly, the algorithm we used in SDSS works here too.

Cosmic ray contaminated pixels satisfy a series of conditions:

- 1. That the candidate bad pixel p not be adjacent to a saturated pixel.
- 2. That p's intensity I exceed the local background by $n\sigma(B)$ ($\sigma(I)$ is the standard deviation of I; $n\approx 6$).
- 3. That no pixel be part of a peak which is sharper than the centre of a star centred in a pixel; i.e.

$$I - c\sigma(I) > c_2\phi(d)\left(\overline{I} + c\sigma(\overline{I})\right)$$

where c and c_2 are constants (\approx 3.0 and 0.6), $\phi(d)$ is the value of the PSF at a distance d from the centre of a star, and \overline{I} is the average of two pixels a distance d away from p.

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These conditions are applied sequentially to the pixel being studied using the four pairs of neighbouring pixels (NS, EW, NW-SE, and NE-SW, $d=1,1,\sqrt{2},\sqrt{2}$).

The candidate cosmic ray must exceed condition 2 for all four pairs of neighbours, and condition 3 for at least one pair.

Cosmic ray contaminated pixels satisfy a series of conditions:

- 1. That the candidate bad pixel p not be adjacent to a saturated pixel.
- 2. That p's intensity l exceed the local background by $n\sigma(B)$ ($\sigma(l)$ is the standard deviation of l; $n \approx 6$).
- 3. That no pixel be part of a peak which is sharper than the centre of a star centred in a pixel; *i.e.*

$$I - c\sigma(I) > c_2\phi(d)\left(\overline{I} + c\sigma(\overline{I})\right)$$

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Once a cosmic ray contaminated pixel is identified, its location is noted and its value is replaced by an interpolation based on the pair of pixels that triggered condition 3.

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We then go through the frame again, looking at pixels adjacent to these cosmic ray events. Processing is identical, except that we set $c_2=0$ for these extra contaminated pixels.

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The only change adopted for thick CCDs was to repeat this last stage an extra couple of times.